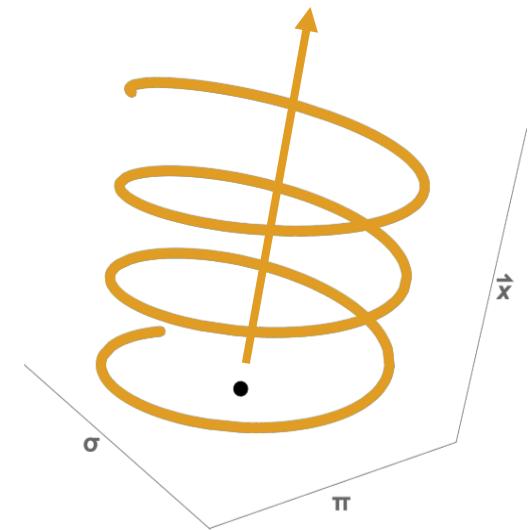


Chiral crossover vs chiral density wave in dense nuclear matter

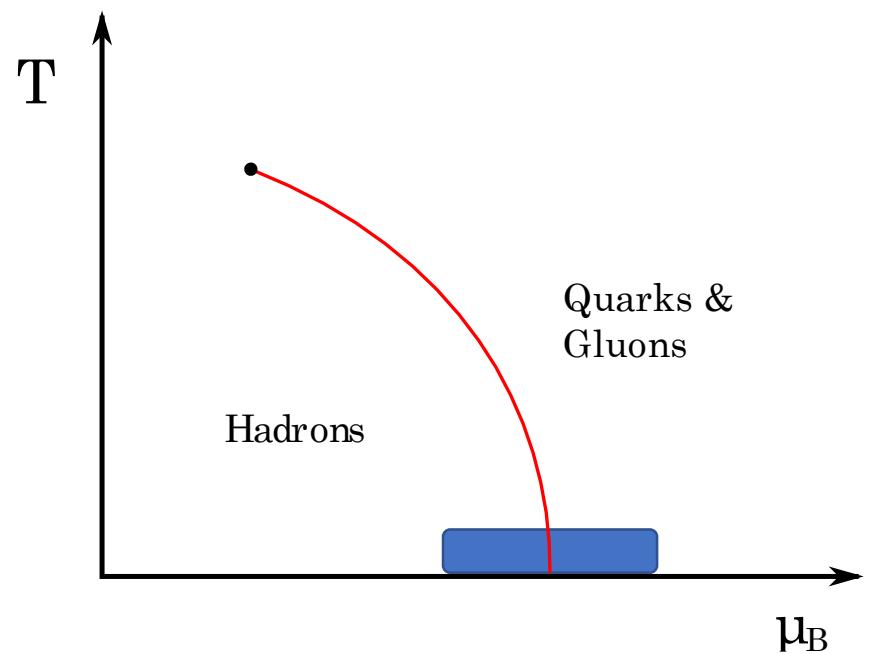
Savvas Pitsinikos

Strong and Electroweak Matter 2024

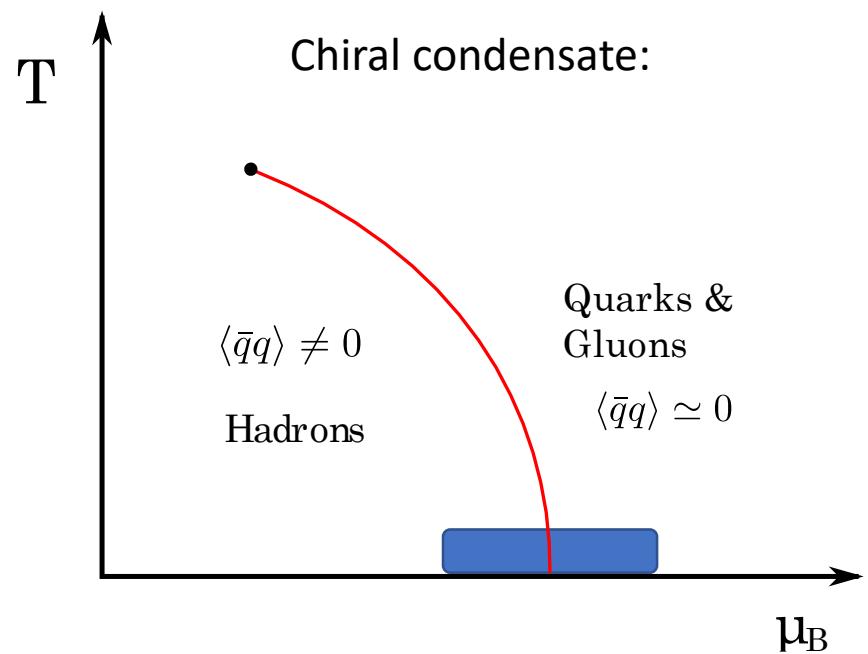
S. Pitsinikos, A. Schmitt, Phys. Rev. D 109, 014024 (2024)



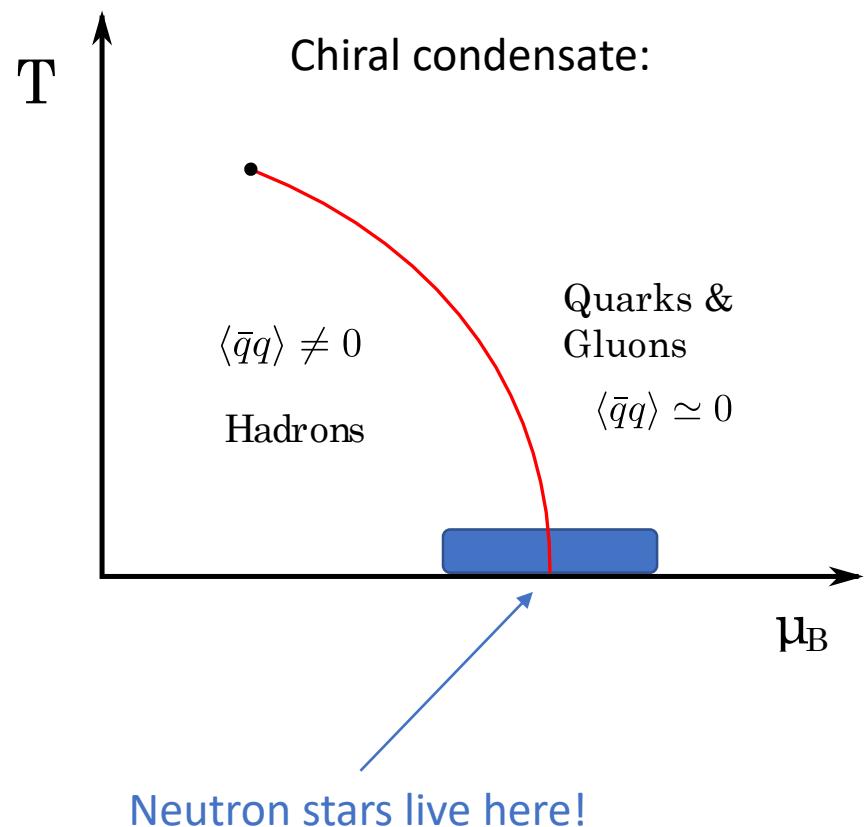
Motivation



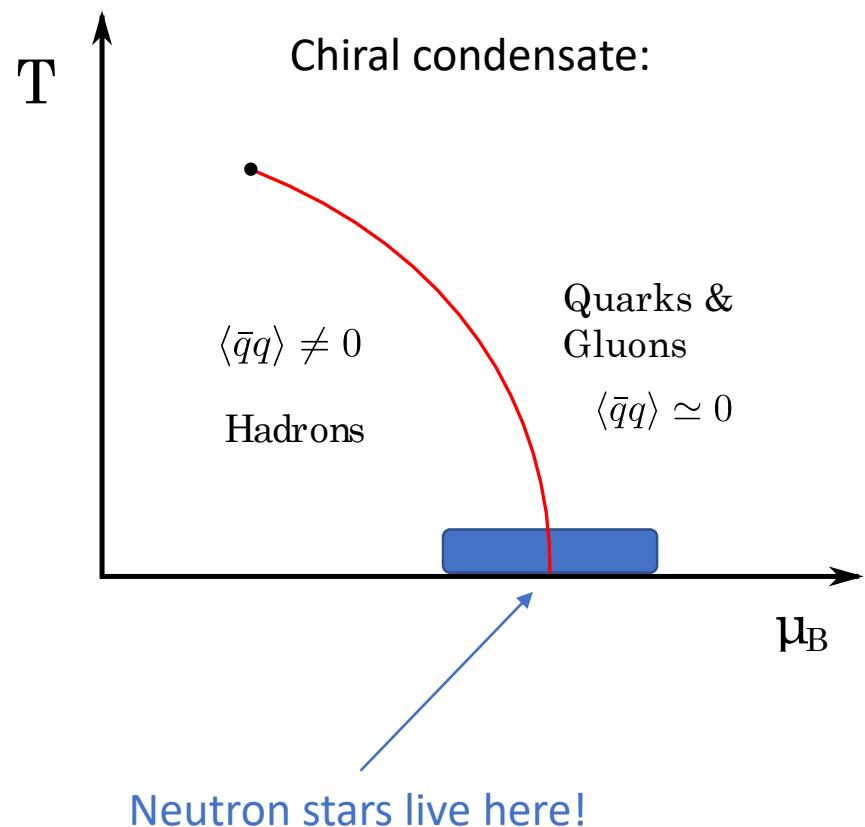
Motivation



Motivation



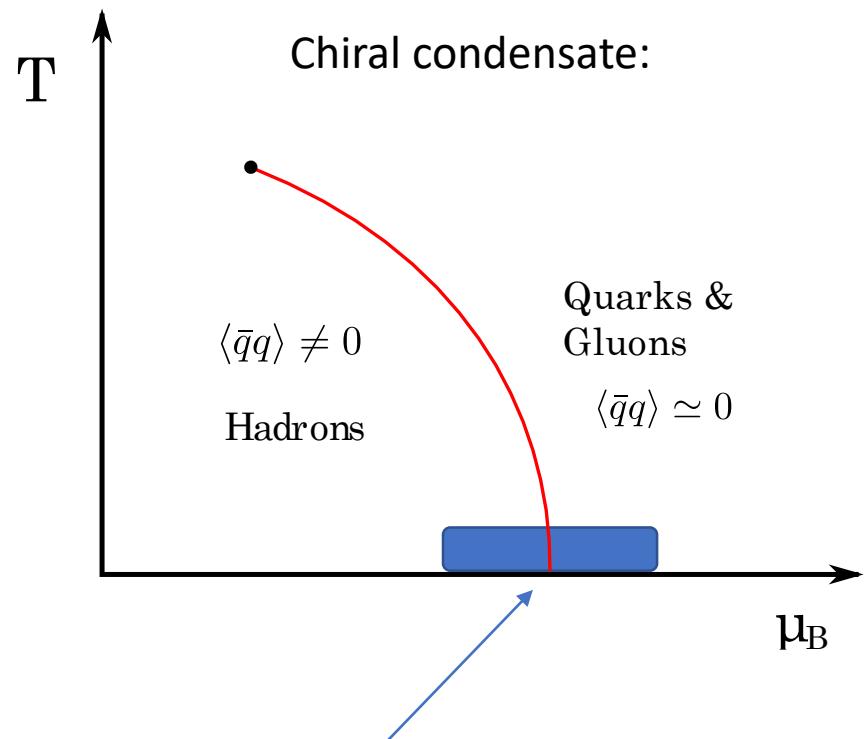
Motivation



What is the
phase structure?

What if the ground
state is not isotropic?

Motivation



Chiral condensate:

$$\langle \bar{q}q \rangle \neq 0$$

Hadrons

Quarks &
Gluons

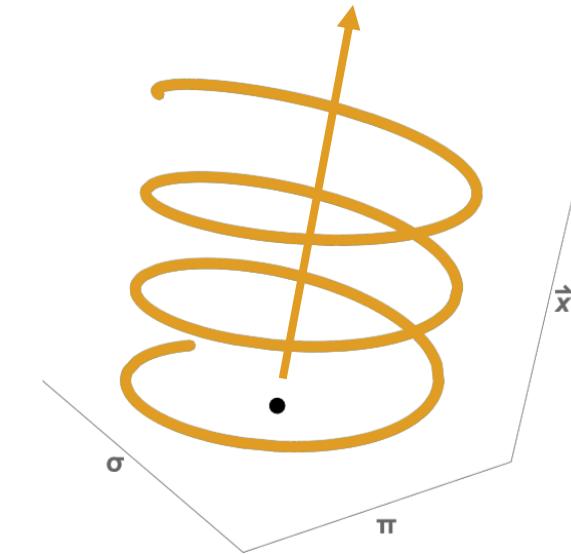
$$\langle \bar{q}q \rangle \simeq 0$$

What is the
phase structure?

What if the ground
state is not isotropic?

Neutron stars live here!

Anisotropic chiral condensate



$$\langle \sigma \rangle = \phi \cos(\vec{q} \cdot \vec{x}), \quad \langle \pi_0 \rangle = \phi \sin(\vec{q} \cdot \vec{x})$$

Chiral Density Wave (CDW)

A. Heinz, F. Giacosa, and D. H. Rischke, Nucl. Phys. A 933, 34 (2015) + references therein
S. Carignano, M. Buballa, PRD, 101(1), 014026 (2020).

Model and novelties

Used to study the chiral phase transition

E. Fraga, R. da Mata, SP, A. Schmitt, Phys. Rev. D 106 7, 074018 (2022)

Andreas Schmitt, Phys. Rev. D 101, 074007

Eduardo S. Fraga, Maurício Hippert, and Andreas Schmitt Phys. Rev. D 99, 014046

Nucleon-meson model:

$$\mathcal{L} = \bar{\psi}(i\gamma^\mu\partial_\mu + \gamma^0\mu)\psi + \frac{1}{2}\partial_\mu\sigma\partial^\mu\sigma + \frac{1}{4}\text{Tr}[\partial_\mu\pi\partial^\mu\pi] - \frac{1}{4}\omega_{\mu\nu}\omega^{\mu\nu} + \frac{m_\omega^2}{2}\omega_\mu\omega^\mu + \frac{d}{4}(\omega_\mu\omega^\mu)^2 - \mathcal{U}(\sigma, \pi)$$
$$-\bar{\psi}[g_\sigma(\sigma + i\gamma^5\pi) + g_\omega\gamma^\mu\omega_\mu]\psi$$

$$\mathcal{U}(\sigma, \pi) = \sum_{n=1}^4 \frac{a_n}{n!} \frac{(\sigma^2 + \pi_a\pi^a - f_\pi^2)^n}{2^n} - \epsilon(\sigma - f_\pi)$$

Model and novelties

Used to study the chiral phase transition

E. Fraga, R. da Mata, SP, A. Schmitt, Phys. Rev. D 106 7, 074018 (2022)

Andreas Schmitt, Phys. Rev. D 101, 074007

Eduardo S. Fraga, Maurício Hippert, and Andreas Schmitt Phys. Rev. D 99, 014046

Nucleon-meson model:

$$\mathcal{L} = \bar{\psi}(i\gamma^\mu\partial_\mu + \gamma^0\mu)\psi + \frac{1}{2}\partial_\mu\sigma\partial^\mu\sigma + \frac{1}{4}\text{Tr}[\partial_\mu\pi\partial^\mu\pi] - \frac{1}{4}\omega_{\mu\nu}\omega^{\mu\nu} + \frac{m_\omega^2}{2}\omega_\mu\omega^\mu + \frac{d}{4}(\omega_\mu\omega^\mu)^2 - \mathcal{U}(\sigma, \pi)$$
$$-\bar{\psi}[g_\sigma(\sigma + i\gamma^5\pi) + g_\omega\gamma^\mu\omega_\mu]\psi$$

$$\mathcal{U}(\sigma, \pi) = \sum_{n=1}^4 \frac{a_n}{n!} \frac{(\sigma^2 + \pi_a\pi^a - f_\pi^2)^n}{2^n} - \epsilon(\sigma - f_\pi)$$

No explicit mass terms for nucleons $M = g_\sigma\phi$,
Chiral symmetry “restoration” possible

Model and novelties

Used to study the chiral phase transition

E. Fraga, R. da Mata, SP, A. Schmitt, Phys. Rev. D 106 7, 074018 (2022)

Andreas Schmitt, Phys. Rev. D 101, 074007

Eduardo S. Fraga, Maurício Hippert, and Andreas Schmitt Phys. Rev. D 99, 014046

Nucleon-meson model:

$$\mathcal{L} = \bar{\psi}(i\gamma^\mu\partial_\mu + \gamma^0\mu)\psi + \frac{1}{2}\partial_\mu\sigma\partial^\mu\sigma + \frac{1}{4}\text{Tr}[\partial_\mu\pi\partial^\mu\pi] - \frac{1}{4}\omega_{\mu\nu}\omega^{\mu\nu} + \frac{m_\omega^2}{2}\omega_\mu\omega^\mu + \frac{d}{4}(\omega_\mu\omega^\mu)^2 - \mathcal{U}(\sigma, \pi)$$
$$-\bar{\psi}[g_\sigma(\sigma + i\gamma^5\pi) + g_\omega\gamma^\mu\omega_\mu]\psi$$

$$\mathcal{U}(\sigma, \pi) = \sum_{n=1}^4 \frac{a_n}{n!} \frac{(\sigma^2 + \pi_a\pi^a - f_\pi^2)^n}{2^n} - \epsilon(\sigma - f_\pi)$$

No explicit mass terms for nucleons $M = g_\sigma\phi$,
Chiral symmetry “restoration” possible

- Mean field free energy and $T = 0$
- Impose isospin symmetry $n_p = n_n$

Model and novelties

Used to study the chiral phase transition

E. Fraga, R. da Mata, SP, A. Schmitt, Phys. Rev. D 106 7, 074018 (2022)

Andreas Schmitt, Phys. Rev. D 101, 074007

Eduardo S. Fraga, Maurício Hippert, and Andreas Schmitt Phys. Rev. D 99, 014046

Nucleon-meson model:

$$\mathcal{L} = \bar{\psi}(i\gamma^\mu\partial_\mu + \gamma^0\mu)\psi + \frac{1}{2}\partial_\mu\sigma\partial^\mu\sigma + \frac{1}{4}\text{Tr}[\partial_\mu\pi\partial^\mu\pi] - \frac{1}{4}\omega_{\mu\nu}\omega^{\mu\nu} + \frac{m_\omega^2}{2}\omega_\mu\omega^\mu + \frac{d}{4}(\omega_\mu\omega^\mu)^2 - \mathcal{U}(\sigma, \pi)$$
$$-\bar{\psi}[g_\sigma(\sigma + i\gamma^5\pi) + g_\omega\gamma^\mu\omega_\mu]\psi$$

$$\mathcal{U}(\sigma, \pi) = \sum_{n=1}^4 \frac{a_n}{n!} \frac{(\sigma^2 + \pi_a\pi^a - f_\pi^2)^n}{2^n} - \epsilon(\sigma - f_\pi)$$

No explicit mass terms for nucleons $M = g_\sigma\phi$,
Chiral symmetry “restoration” possible

- Mean field free energy and $T = 0$
- Impose isospin symmetry $n_p = n_n$

Similar in spirit approaches:

- Inhomogeneous phases extensively studied in quark models

S. Carignano, M. Buballa, B.J. Schaefer, Phys. Rev. D 90, 014033 (2014)

- Nucleon models ignore Dirac sea

A. Heinz, F. Giacosa, and D. H. Rischke, Nucl. Phys. A 933, 34 (2015)

Dirac sea and renormalization

Dirac sea and renormalization

Why?

Dirac sea and renormalization

Why?

Significant contribution with qualitative effects
for the chiral transition

Brandes L., Kaiser N. & Weise W. Eur. Phys. J. A 57, 243 (2021)

1st order → Crossover

Dirac sea and renormalization

Why?

Baryonic vacuum contribution:

$$P_{\text{vac}}(\phi, q) = \sum_{n,p} \int \frac{d^3 \vec{k}}{(2\pi)^3} E_k(\phi, q)$$

divergent!

Significant contribution with qualitative effects
for the chiral transition

Brandes L., Kaiser N. & Weise W. Eur. Phys. J. A 57, 243 (2021)
1st order → Crossover

Dirac sea and renormalization

Why?

Baryonic vacuum contribution:

$$P_{\text{vac}}(\phi, q) = \sum_{n,p} \int \frac{d^3 \vec{k}}{(2\pi)^3} E_k(\phi, q)$$

divergent!

Significant contribution with qualitative effects
for the chiral transition

Brandes L., Kaiser N. & Weise W. Eur. Phys. J. A 57, 243 (2021)
1st order → Crossover

Proper time regularization Λ + renormalization at scale ℓ

Dirac sea and renormalization

Why?

Baryonic vacuum contribution:

$$P_{\text{vac}}(\phi, q) = \sum_{n,p} \int \frac{d^3 \vec{k}}{(2\pi)^3} E_k(\phi, q)$$

divergent!

Significant contribution with qualitative effects
for the chiral transition

Brandes L., Kaiser N. & Weise W. Eur. Phys. J. A 57, 243 (2021)
1st order → Crossover

Proper time regularization Λ + renormalization at scale ℓ

Select $\ell = \sqrt{m_N^2 + (2q)^2}$

Eduardo S. Fraga, Letícia F. Palhares, and Túlio E. Restrepo
Phys. Rev. D 108, 034026 (2023)
Aleksi Kurkela, Paul Romatschke, and Aleksi Vuorinen
Phys. Rev. D 81, 105021 (2010)

Dirac sea and renormalization

Why?

Baryonic vacuum contribution:

$$P_{\text{vac}}(\phi, q) = \sum_{n,p} \int \frac{d^3 \vec{k}}{(2\pi)^3} E_k(\phi, q)$$

divergent!

Significant contribution with qualitative effects
for the chiral transition

Brandes L., Kaiser N. & Weise W. Eur. Phys. J. A 57, 243 (2021)
1st order → Crossover

Proper time regularization Λ + renormalization at scale ℓ

Select $\ell = \sqrt{m_N^2 + (2q)^2}$

Eduardo S. Fraga, Letícia F. Palhares, and Túlio E. Restrepo
Phys. Rev. D 108, 034026 (2023)
Aleksi Kurkela, Paul Romatschke, and Aleksi Vuorinen
Phys. Rev. D 81, 105021 (2010)

So that free energy is bounded for large q :

$$\Omega = g_\sigma^2 \frac{q^2 \phi^2}{2\pi^2} \left(2 + \frac{4\pi^2}{g_\sigma^2} - \ln \frac{4q^2}{\ell^2} \right) + \mathcal{O}(q^0)$$

Fixing the model parameters

Free parameters \leftrightarrow properties of symmetric nuclear matter at saturation:

- Binding energy: $E_B = -16.3 \text{ MeV}$
- Density: $n_0 = 0.15 \text{ fm}^{-3}$
- Incompressibility: $K \approx (200 - 300) \text{ MeV}$
- Symmetry energy: $S \approx (30 - 34) \text{ MeV}$
- Slope of symmetry energy: $L \approx (40 - 140) \text{ MeV}$ (PREX), Phys. Rev. Lett. 126, 172502 (2021)
- Effective nucleon mass: $M_0 \approx (0.7 - 0.8)m_N$

Fixing the model parameters

Free parameters \leftrightarrow properties of symmetric nuclear matter at saturation:

- Binding energy: $E_B = -16.3 \text{ MeV}$
- Density: $n_0 = 0.15 \text{ fm}^{-3}$
- Incompressibility: $K \approx (200 - 300) \text{ MeV}$
- Symmetry energy: $S \approx (30 - 34) \text{ MeV}$
- Slope of symmetry energy: $L \approx (40 - 140) \text{ MeV}$ (PREX), Phys. Rev. Lett. 126, 172502 (2021)
- Effective nucleon mass: $M_0 \approx (0.7 - 0.8)m_N$

We mainly explore the
 M_0, L
parameter space

Fixing the model parameters

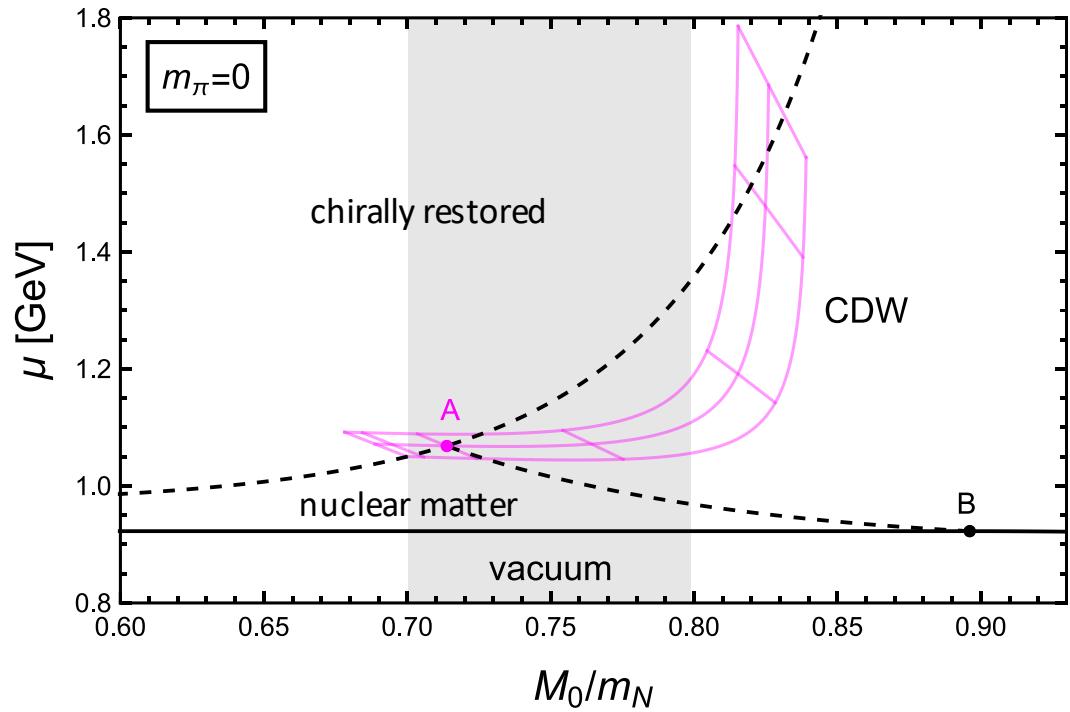
Free parameters \leftrightarrow properties of symmetric nuclear matter at saturation:

- Binding energy: $E_B = -16.3 \text{ MeV}$
- Density: $n_0 = 0.15 \text{ fm}^{-3}$
- Incompressibility: $K \approx (200 - 300) \text{ MeV}$
- Symmetry energy: $S \approx (30 - 34) \text{ MeV}$
- Slope of symmetry energy: $L \approx (40 - 140) \text{ MeV}$ (PREX), Phys. Rev. Lett. 126, 172502 (2021)
- Effective nucleon mass: $M_0 \approx (0.7 - 0.8)m_N$

We mainly explore the
 M_0, L
parameter space

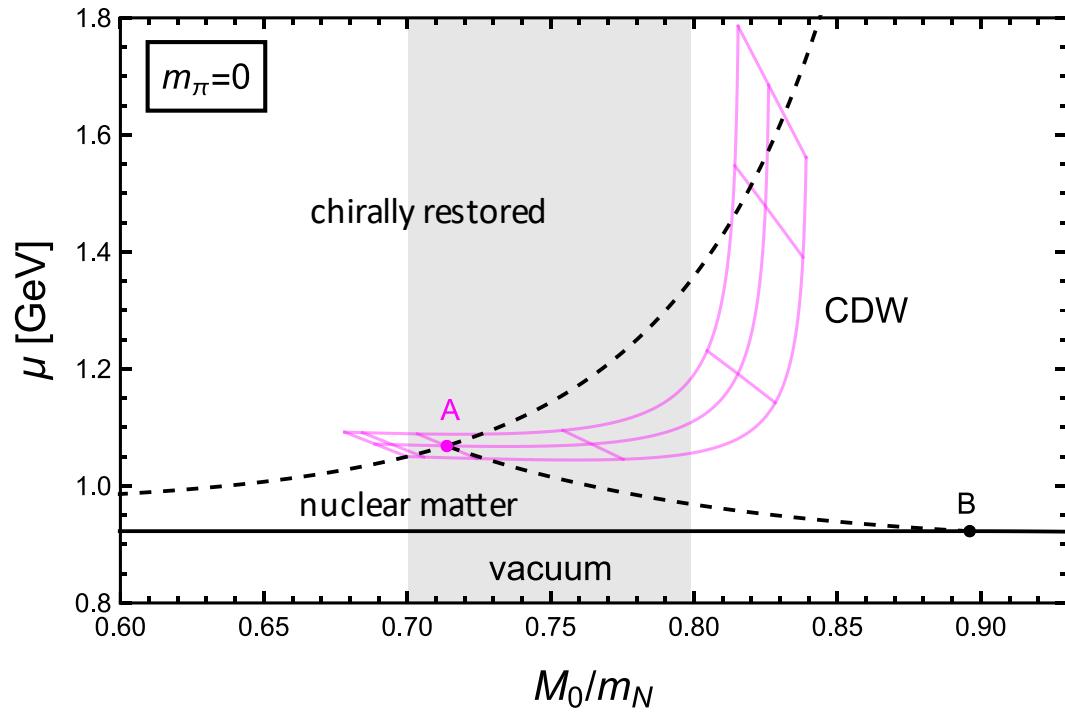
Direct relation to nuclear matter phenomenology!

Results - CDW

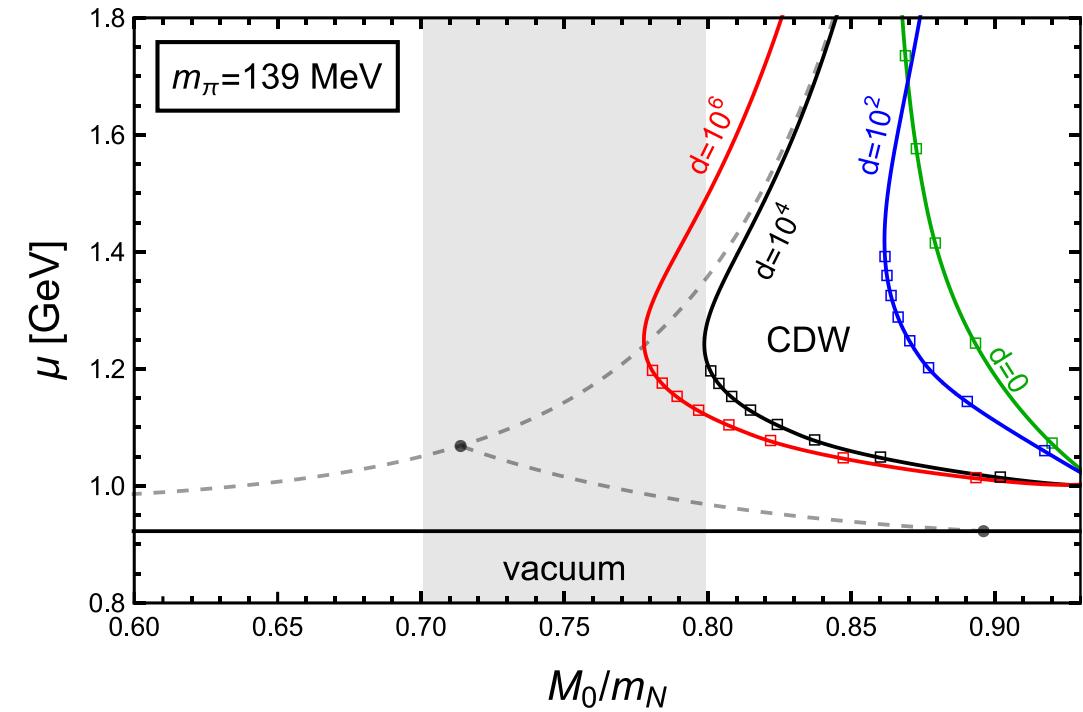


CDW replaces the chiral phase transition

Results - CDW

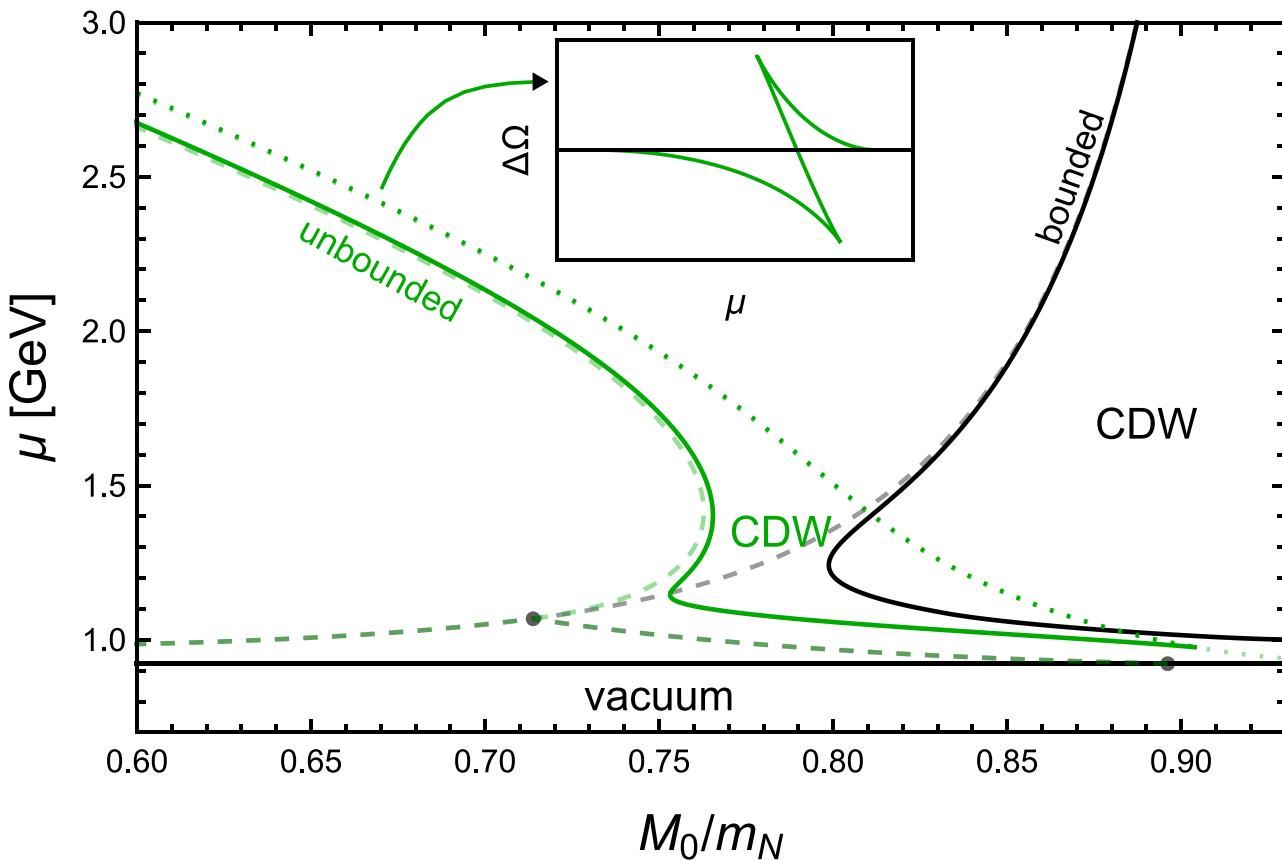


CDW replaces the chiral phase transition



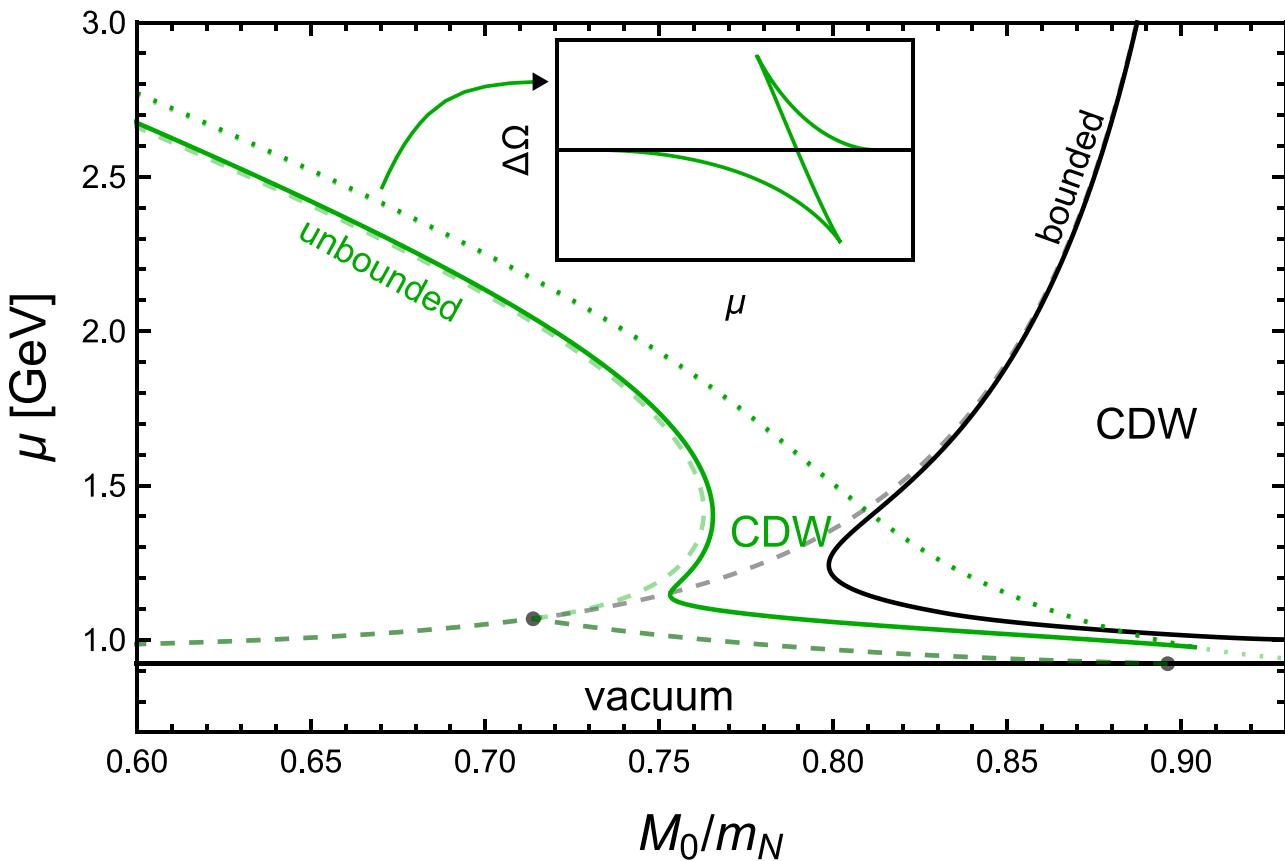
Explicit chiral symmetry breaking shifts the CDW

Results - CDW

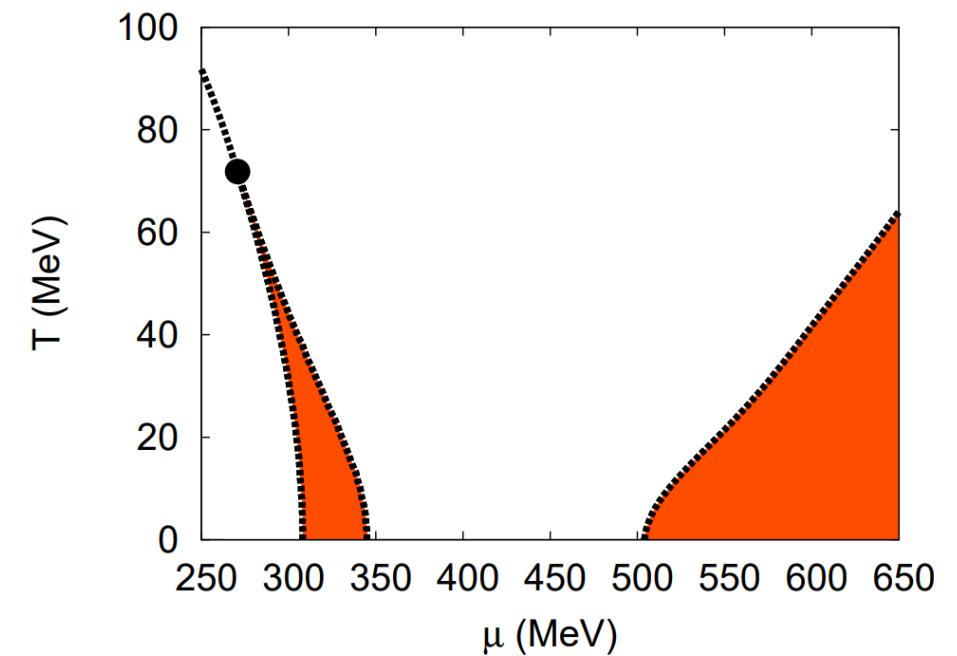


Without the q -dependent renormalization scale
the unphysical island appears

Results - CDW



Without the q -dependent renormalization scale
the unphysical island appears

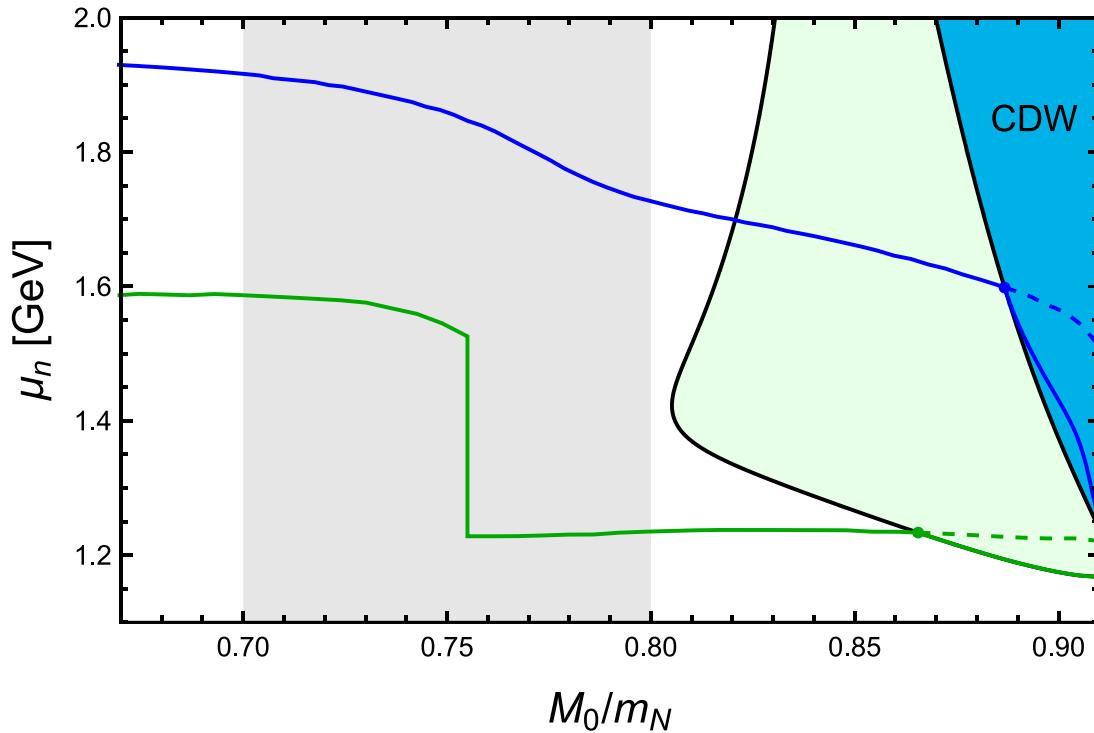


CDW in Neutron Stars?

Preliminary

Poster: "Chiral Density Wave under neutron star conditions"

O. Papadopoulos and A. Schmitt – Work in progress

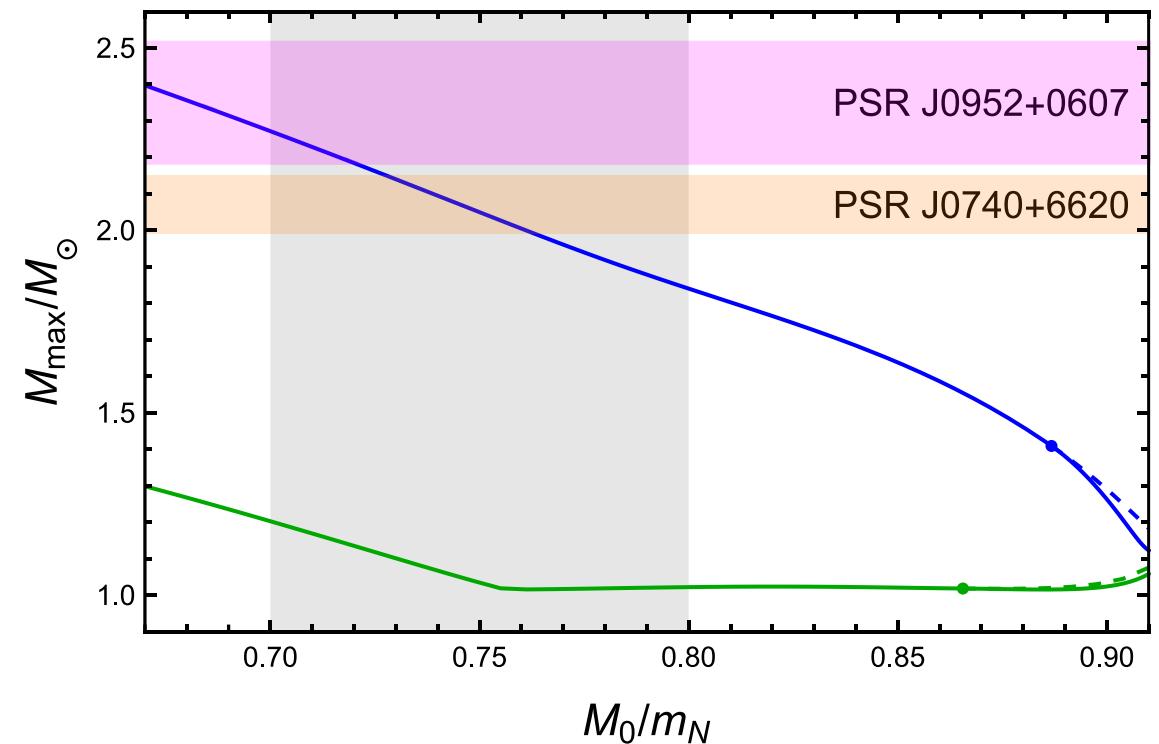
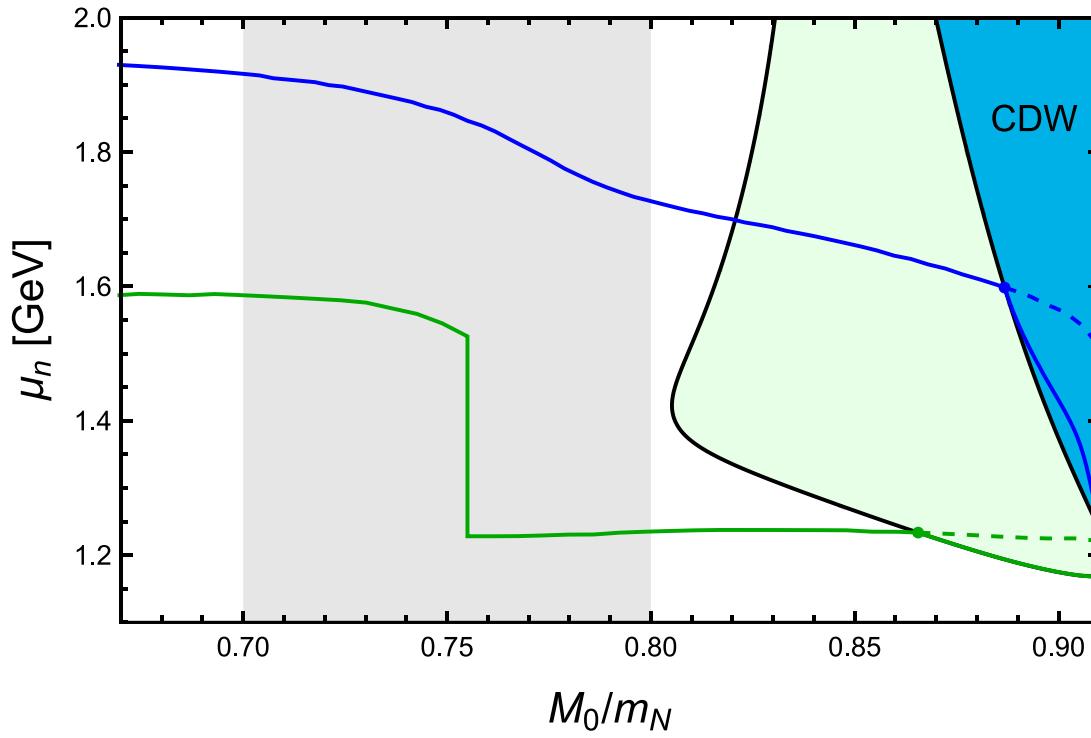


CDW in Neutron Stars?

Preliminary

Poster: "Chiral Density Wave under neutron star conditions"

O. Papadopoulos and A. Schmitt – Work in progress

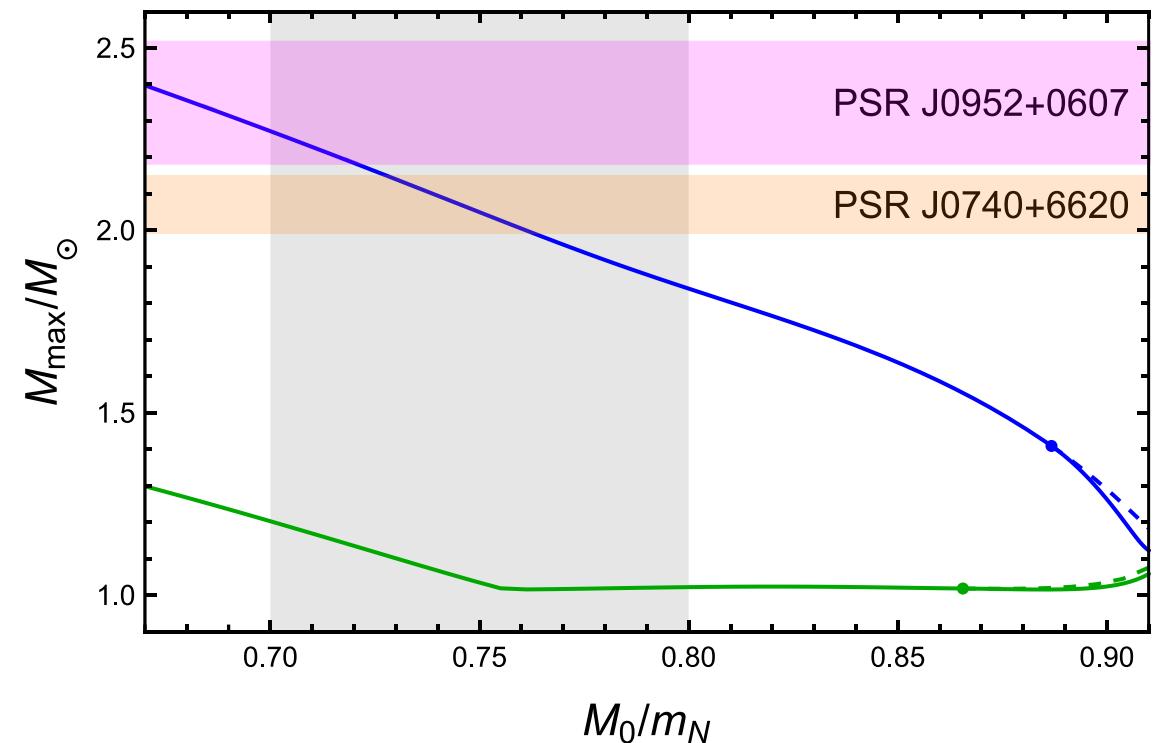
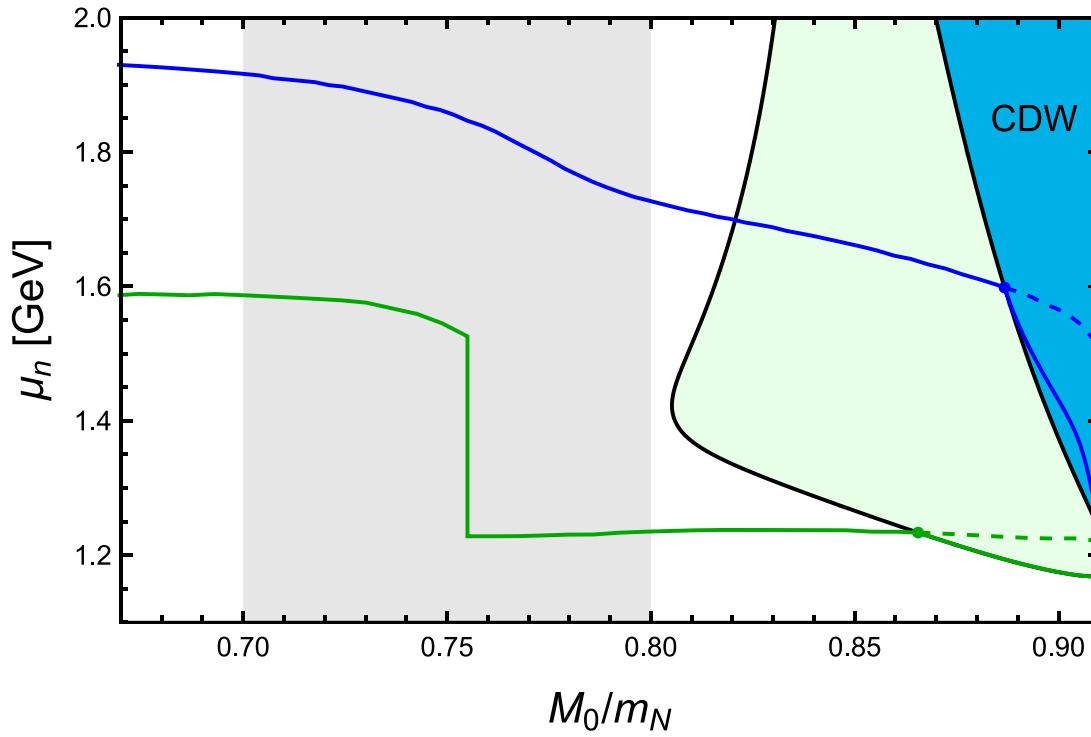


CDW in Neutron Stars?

Preliminary

Poster: "Chiral Density Wave under neutron star conditions"

O. Papadopoulos and A. Schmitt – Work in progress



Existence of heavy compact stars excludes CDW

Summary and Outlook

Summary and Outlook

- Setup a nucleonic model for the CDW

Summary and Outlook

- Setup a nucleonic model for the CDW
- Included the Dirac sea

Summary and Outlook

- Setup a nucleonic model for the CDW
- Included the Dirac sea
- Used a renormalization scheme to cure previously observed unphysical behavior

Summary and Outlook

- Setup a nucleonic model for the CDW
- Included the Dirac sea
- Used a renormalization scheme to cure previously observed unphysical behavior
- Mapped the region in the parameter space where CDW is preferred

Summary and Outlook

- Setup a nucleonic model for the CDW
- Included the Dirac sea
- Used a renormalization scheme to cure previously observed unphysical behavior
- Mapped the region in the parameter space where CDW is preferred

Outlook

- Beta equilibrium, charge neutrality O. Papadopoulos and A. Schmitt – In preparation
- Alternative models to tackle model dependent effects A. Heinz, F. Giacosa, D. H. Rischke, Nucl. Phys. A, 933, 34-42, (2015)
- Transport properties may be significantly modified With implications for:
M. Alford, A. Harutyunyan, A. Sedrakian,
Phys. Rev. D 108, 083019 (2023)

Summary and Outlook

- Setup a nucleonic model for the CDW
- Included the Dirac sea
- Used a renormalization scheme to cure previously observed unphysical behavior
- Mapped the region in the parameter space where CDW is preferred

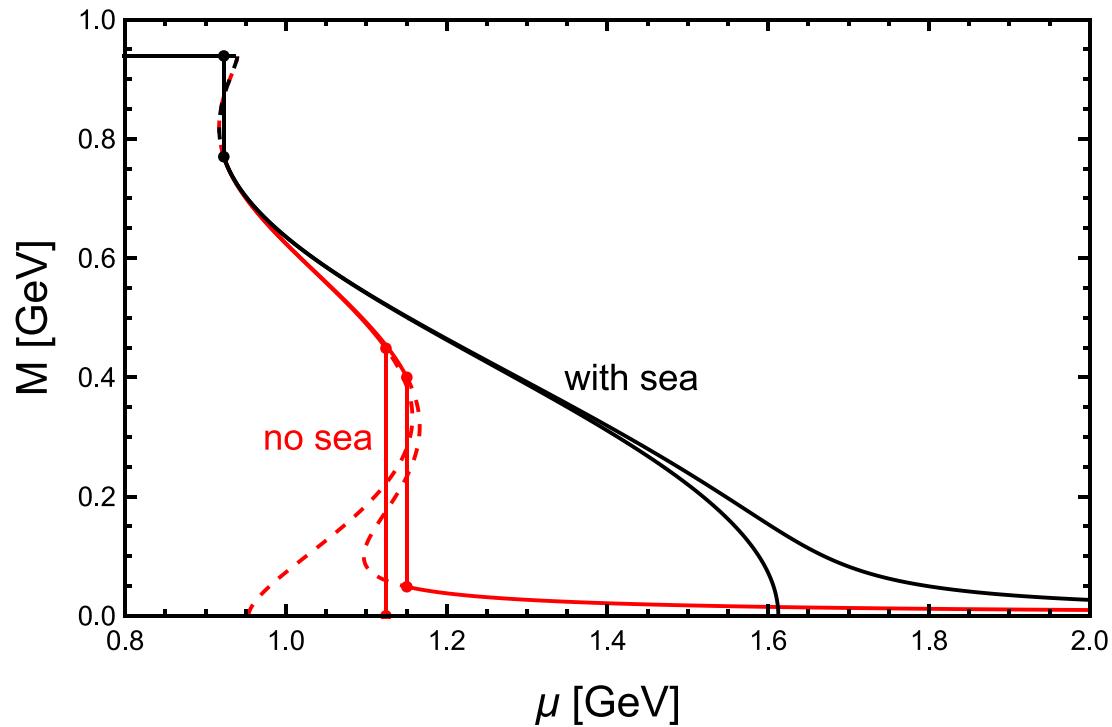
Outlook

- Beta equilibrium, charge neutrality O. Papadopoulos and A. Schmitt – In preparation
- Alternative models to tackle model dependent effects A. Heinz, F. Giacosa, D. H. Rischke, Nucl. Phys. A, 933, 34-42, (2015)
- Transport properties may be significantly modified With implications for:
M. Alford, A. Harutyunyan, A. Sedrakian,
Phys. Rev. D 108, 083019 (2023)

Thank you!

Extras

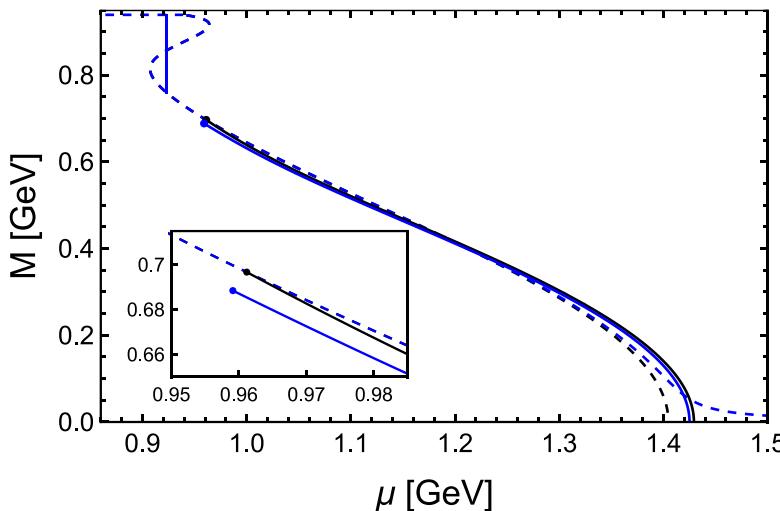
Results - Isotropic



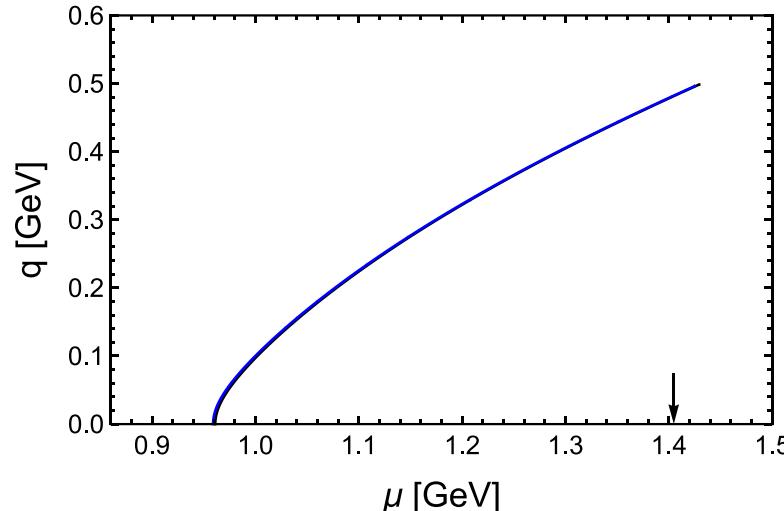
The Dirac sea's effect is important even for homogeneous phases

Brandes L., Kaiser N. & Weise W. Eur. Phys. J. A 57, 243 (2021)

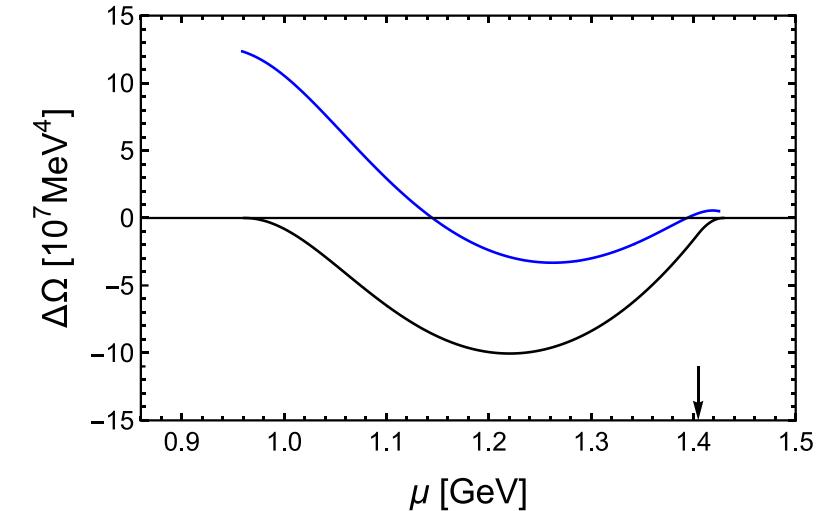
Results - CDW



CDW appears despite the crossover



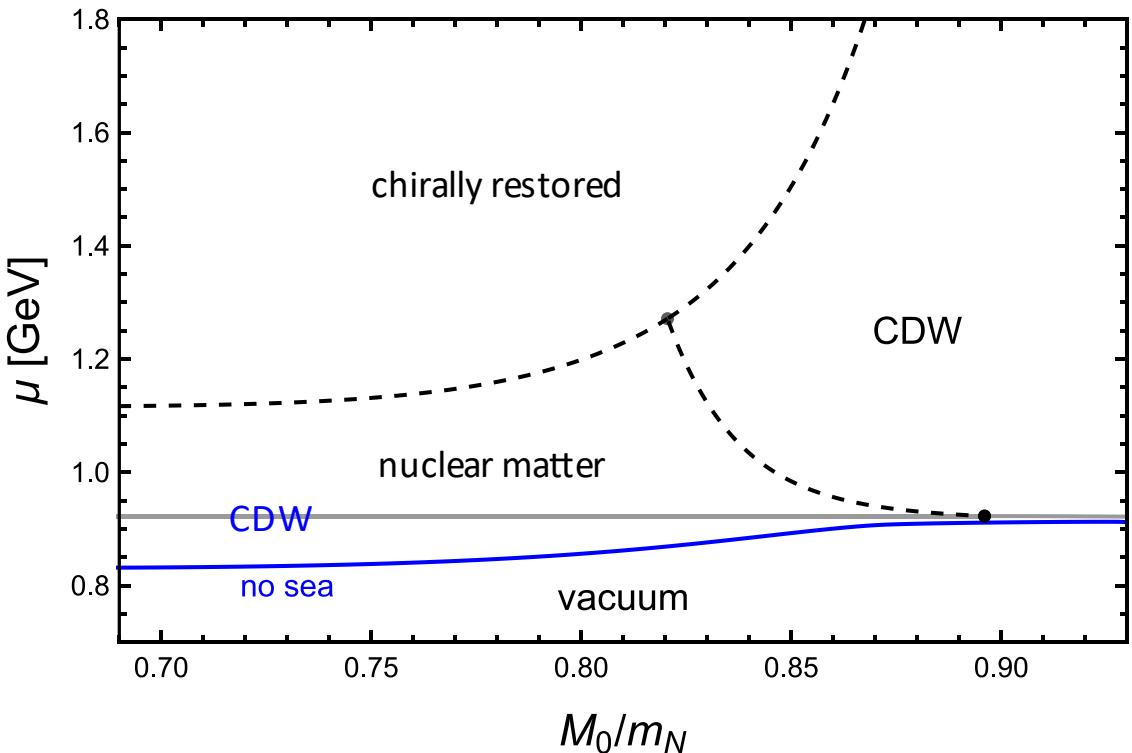
Small effect of the pion mass



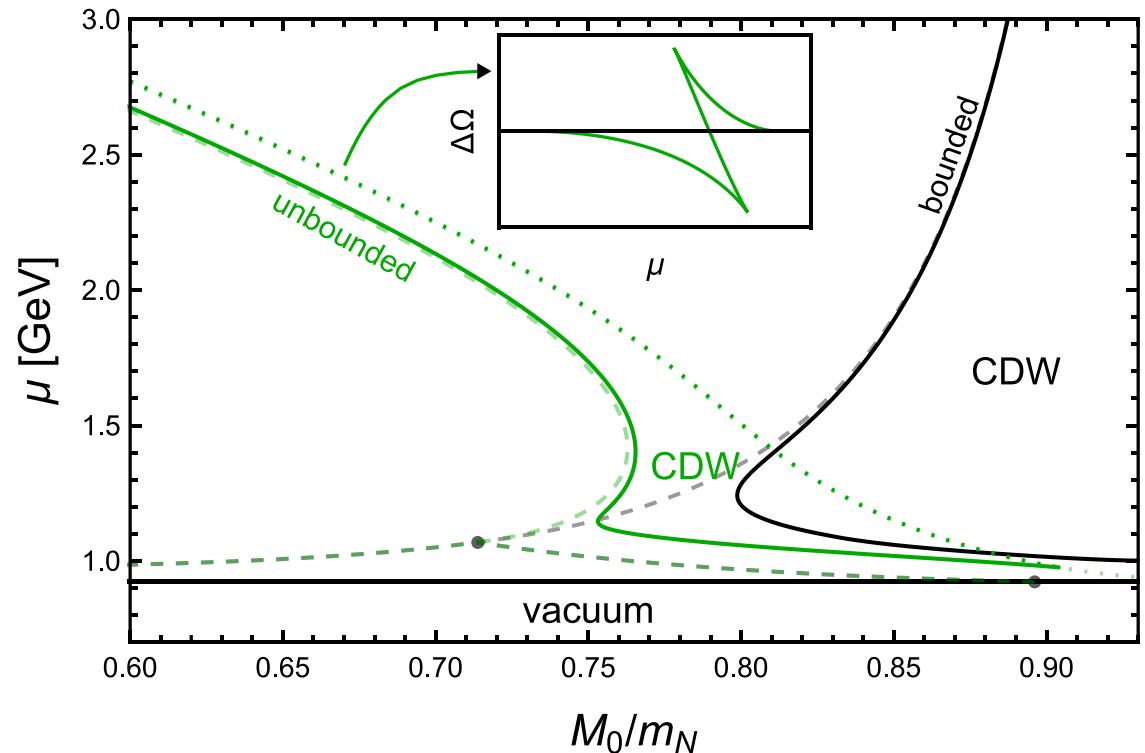
CDW preferred when $\Delta\Omega < 0$

Chiral limit – Physical pion mass

Results - CDW



Without the sea the CDW is always preferred



Without the q -dependent renormalization scale
the unphysical island appears